

# Sorption and Transport Properties of Propane and Perfluoropropane in Poly(dimethylsiloxane) and Poly(1-trimethylsilyl-1-propyne)

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**ABSTRACT:** The effect of pressure on solubility and the influence of temperature on solubility, permeability, and diffusivity of C<sub>3</sub>F<sub>8</sub> and its hydrocarbon analogue, C<sub>3</sub>H<sub>8</sub>, are reported in rubbery poly(dimethylsiloxane) (PDMS) and glassy poly(1-trimethylsilyl-1-propyne) (PTMSP). The solubility of C<sub>3</sub>F<sub>8</sub> is lower than that of C<sub>3</sub>H<sub>8</sub> in both polymers at all temperatures and pressures investigated. The isosteric enthalpy of mixing C<sub>3</sub>F<sub>8</sub> with PDMS and PTMSP is higher than that of C<sub>3</sub>H<sub>8</sub> due to less favorable polymer–fluorocarbon interactions in the case of C<sub>3</sub>F<sub>8</sub>, and it decreases with increasing C<sub>3</sub>F<sub>8</sub> concentration. Assuming a coordination number of 10, the energy associated with mixing C<sub>3</sub>F<sub>8</sub> molecules and PDMS segments is 4.5 kJ/mol more than that required to mix C<sub>3</sub>H<sub>8</sub> molecules with PDMS segments, in the limit of infinite dilution. The isobaric activation energy of permeation ( $E_P$ ) for C<sub>3</sub>F<sub>8</sub> is positive for both polymers, and that for C<sub>3</sub>H<sub>8</sub> is negative in both polymers. This result is particularly interesting for PTMSP since all previous studies of activation energy of gas permeation in PTMSP report values that are near zero or negative; this study provides the first report of a positive  $E_P$  value in PTMSP. In PDMS, differences in both activation energy of diffusion ( $E_D$ ) and enthalpy change on sorption contribute significantly to the difference in  $E_P$  values of C<sub>3</sub>H<sub>8</sub> and C<sub>3</sub>F<sub>8</sub>. For PTMSP, the difference in  $E_P$  values for C<sub>3</sub>F<sub>8</sub> and C<sub>3</sub>H<sub>8</sub> stems mainly from a substantially larger  $E_D$  value for C<sub>3</sub>F<sub>8</sub> than for C<sub>3</sub>H<sub>8</sub>.

## Introduction

Perfluorocarbons (PFCs) are industrially important compounds with very long atmospheric lifetimes (3000–5000 years) and enormous global warming potentials (5000–10 000 times greater than that of CO<sub>2</sub>).<sup>1</sup> Because of the potentially deleterious effect of these materials on the environment, U.S. federal regulations have restricted their use to “high-performance, precision-engineered applications only, where reasonable efforts have been made to ascertain that other alternatives are not technically feasible due to performance or safety requirements.”<sup>2</sup>

The separation of PFCs from permanent gases using membranes has been the subject of several investigations.<sup>3–10</sup> Relative to their hydrocarbon analogues, PFCs exhibit low solubility in hydrocarbon-based polymers such as poly(dimethylsiloxane) (PDMS),<sup>7,11</sup> poly(1-trimethylsilyl-1-propyne) (PTMSP),<sup>5,7</sup> and polyethylene.<sup>11</sup> This behavior has been ascribed to less favorable interactions between fluorinated penetrants and hydrocarbon-based polymers than between hydrocarbon penetrants and hydrocarbon-based polymers.<sup>7,9,10</sup> Similar observations of low fluorocarbon solubility in hydrocarbon liquids have also been reported.<sup>12–14</sup> For example, the mole fractions of C<sub>3</sub>F<sub>8</sub> and C<sub>3</sub>H<sub>8</sub> dissolved in cyclohexane at 25 °C and 1 atm are 0.006 532 and 0.095 578, respectively;<sup>14</sup> that is, C<sub>3</sub>F<sub>8</sub> is almost 15 times less soluble in cyclohexane than C<sub>3</sub>H<sub>8</sub>.

This report provides quantitative, experimental evidence of the less favorable interactions between fluorocarbon penetrants and hydrocarbon-based polymers that influence the energetics of gas sorption and transport. Perfluoropropane was selected as a model penetrant, and its sorption, diffusion, and permeation properties are compared with those of its hydrocarbon analogue, propane, in two very different hydrocarbon-based polymers, PDMS and PTMSP. PDMS is a rubbery polymer (its glass transition temperature,  $T_g$ , is –123 °C).<sup>15</sup> As such, it presents a mobile, liquidlike environment to penetrant molecules. PTMSP, on the other hand, is a stiff chain, glassy polymer ( $T_g > 250$  °C) exhibiting very poor chain packing in the solid state.<sup>16,17</sup> It is the most permeable polymer known, and it has the lowest density and highest fractional free volume of all known hydrocarbon-based polymers.<sup>9</sup> Permeability coefficients of N<sub>2</sub> and H<sub>2</sub> in these polymers are also provided because separation of PFCs from mixtures with these permanent gases has been the focus of industrial interest.<sup>18–26</sup>

## Background

**Permeability.** The permeability of a polymer film to a pure penetrant is given by<sup>27</sup>

$$P = \frac{Nl}{p_2 - p_1} \quad (1)$$

where  $P$  is the permeability coefficient,  $N$  is the steady-state gas flux through the polymer membrane,  $l$  is film thickness,  $p_2$  is the feed or upstream pressure, and  $p_1$  is the permeate or downstream pressure. Penetrant

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transport through a polymer film is commonly described by a three-step solution-diffusion mechanism.<sup>27</sup> According to this mechanism, when penetrant flux obeys Fick's law

$$P = \frac{C_2 - C_1}{p_2 - p_1} D \quad (2)$$

where  $D$  is the effective concentration-averaged diffusion coefficient, and  $C_2$  and  $C_1$  are the penetrant concentrations at the upstream and downstream faces of the membrane. If the downstream pressure is negligible compared to the upstream pressure, eq 2 can be simplified:<sup>27</sup>

$$P = SD \quad (3)$$

where the solubility coefficient,  $S$ , is defined as follows:

$$S = \frac{C}{p} \quad (4)$$

In eq 3,  $S$  should be evaluated at the upstream conditions. The effective concentration-averaged diffusivity is defined by

$$D = \int_{C_1}^{C_2} \frac{D_{\text{loc}}}{1 - \omega} dC = \int_{C_1}^{C_2} D_{\text{eff}} dC \quad (5)$$

where  $D_{\text{loc}}$  is the local concentration-dependent diffusion coefficient,  $\omega$  is the penetrant mass fraction in the polymer at concentration  $C$ , and  $D_{\text{eff}}$  is the so-called local effective diffusion coefficient.

**Selectivity.** The ideal selectivity,  $\alpha_{A/B}$ , of component A over B is a measure of the potential separation ability of the membrane material. The ideal selectivity can be written as the ratio of the pure gas permeabilities:<sup>27</sup>

$$\alpha_{A/B} \equiv \frac{P_A}{P_B} \quad (6)$$

or

$$\alpha_{A/B} = \frac{S_A D_A}{S_B D_B} \quad (7)$$

where the first term on the right-hand side is the solubility selectivity and the second is the diffusivity selectivity. In addition to operating conditions (i.e., temperature, pressure, and gas composition), penetrant solubility depends on condensability and polymer-penetrant interactions.<sup>27</sup> In the absence of specific interactions (e.g., hydrogen bonding), the first effect is dominant, and solubility increases as penetrant condensability, characterized by critical temperature, normal boiling point, or Lennard-Jones force constant, increases.<sup>27</sup> Thus, solubility selectivity increases as the difference in condensability between two penetrants in a mixture increases. Often, larger penetrants are more condensable and therefore more soluble than smaller penetrants. However, larger penetrants have lower diffusion coefficients than smaller penetrants, and as the difference in penetrant size increases, so does the diffusivity selectivity.<sup>27</sup> Thus, a tradeoff often exists between solubility selectivity and diffusivity selectivity, with the overall selectivity depending on the relative magnitudes of these two terms.

**Solubility.** The sorption of sparingly soluble gases in rubbery polymers is qualitatively similar to the sorption of gases in low molecular weight liquids, and the gas concentration in the polymer,  $C$ , often obeys Henry's law:<sup>27</sup>

$$C = k_D p \quad (8)$$

where  $k_D$  is the Henry's law constant and  $p$  is the gas pressure in contact with the polymer. The uptake of more soluble vapors in a cross-linked rubbery polymer is frequently described using the Flory-Rehner expression:<sup>28</sup>

$$\ln a = \ln \phi_2 + (1 - \phi_2) + \chi(1 - \phi_2)^2 + V_2 \left( \frac{v_e}{V_0} \right) \left[ (1 - \phi_2)^{1/3} - \left( \frac{1 - \phi_2}{2} \right) \right] \quad (9)$$

where  $a$  is penetrant activity in the vapor phase,  $\phi_2$  is the volume fraction of sorbed penetrant,  $\chi$  is the Flory-Huggins interaction parameter,  $V_2$  is the penetrant molar volume, and  $v_e/V_0$  is the effective number of cross-links per unit volume of penetrant-free polymer (expressed in moles of cross-links per unit volume of penetrant-free polymer). In this study, penetrant activity is set equal to the relative pressure,  $p/p_{\text{sat}}$ , where  $p_{\text{sat}}$  is the saturation vapor pressure of the penetrant. The volume fraction of sorbed penetrant,  $\phi_2$ , is calculated from the equilibrium penetrant concentration in the polymer,  $C$ , as follows:

$$\phi_2 = \left[ 1 + \frac{22414}{C \bar{V}_2} \right]^{-1} \quad (10)$$

where  $\bar{V}_2$  is the penetrant partial molar volume and is estimated as described previously.<sup>10</sup> In this equation,  $C$  and  $\bar{V}_2$  have units of  $\text{cm}^3$  (STP)/ $\text{cm}^3$  polymer and  $\text{cm}^3/\text{mol}$ , respectively. 22414 is a conversion factor ( $\text{cm}^3$  (STP)/mol).

Sorption isotherms for gases in glassy polymers are usually concave to the pressure axis at low pressures and linear at higher pressures.<sup>27</sup> Such isotherms are often described using the dual-mode sorption model.<sup>29</sup> In this model, penetrant molecules are viewed as being partitioned into two populations which are in dynamic equilibrium with each other: (i) penetrant molecules sorbed by a dissolution mechanism in the dense polymer matrix (Henry's law population) and (ii) penetrant molecules filling unrelaxed, molecular-scale gaps (microvoids) frozen into the glassy state (Langmuir population).<sup>29</sup> The dual-mode model is expressed analytically as a sum of these two contributions to penetrant sorption:

$$C = k_D p + \frac{C'_H b p}{1 + b p} \quad (11)$$

where  $C$  is the total concentration of penetrant in the polymer,  $k_D$  is the Henry's law constant,  $C'_H$  is the hole saturation constant or Langmuir sorption capacity parameter, and  $b$  is the Langmuir affinity parameter.

**Temperature Dependence.** The temperature dependence of permeability, diffusivity, and solubility at temperatures far removed from polymer thermal transitions is described as follows:<sup>27</sup>

$$P = P_0 \exp\left(-\frac{E_p}{RT}\right) \quad (12)$$

$$D = D_0 \exp\left(-\frac{E_D}{RT}\right) \quad (13)$$

$$S = S_0 \exp\left(-\frac{\Delta H_s}{RT}\right) \quad (14)$$

where  $P_0$ ,  $D_0$ , and  $S_0$  are preexponential constants,  $E_p$  is the activation energy of permeation,  $E_D$  is the activation energy of diffusion, and  $\Delta H_s$  is the enthalpy of sorption. Because permeability is the product of solubility and diffusivity (eq 3), the activation energies of permeation and diffusion and the enthalpy of sorption are related:

$$E_p = E_D + \Delta H_s \quad (15)$$

Equation 15 is a consequence of eqs 3, 12, 13, and 14 and is, therefore, subject to the assumptions inherent in these equations. For example, eq 15 does not hold if the downstream pressure cannot be neglected in comparison to the upstream pressure due to the assumptions underlying eq 3. Also, if the penetrant transport properties (i.e.,  $P$ ,  $D$ , and  $S$ ) are functions of concentration, eq 15 is expected to be a simplified form of a more general model.<sup>30</sup> Nevertheless, eqs 12–15 are the standard model for describing the temperature dependence of gas solubility, diffusivity, and permeability in polymers, and we will use them in this study.

## Experimental Section

**Materials.** PDMS. Poly(dimethylsiloxane) (PDMS) composite membranes were used for pure gas permeation experiments. These membranes, composed of a filler-free PDMS film on a highly microporous support, were kindly provided by Dr. Ingo Pinnau of Membrane Technology and Research, Inc. (Menlo Park, CA). The poly(dimethylsiloxane) was from Wacker Silicones Corp. (Adrian, MI) and was cross-linked at 100 °C using a proprietary cross-linker/catalyst system supplied by them.

A dense filler-free PDMS film of thickness approximately 250  $\mu\text{m}$  was used for the sorption measurements. Cross-linking was achieved using the same method described above. The cross-link density of this film was estimated to be  $7.8 \times 10^{-5}$  mol/cm<sup>3</sup>.<sup>31</sup> Since the film for the permeation measurements was cross-linked under the same conditions, it should have a similar cross-link density to that of the dense film.

**PTMSP.** PTMSP was kindly provided by Permea, Inc. (St. Louis, MO). Isotropic poly(1-trimethylsilyl-1-propyne) films, approximately 50  $\mu\text{m}$  thick, were prepared from a 2 wt % solution of the polymer in toluene as described previously.<sup>32</sup> After casting and drying, the samples were stored in liquid methanol at ambient conditions to mitigate physical aging. The films were removed from methanol and dried at ambient conditions for 24 h before using them for experiments. These films were utilized for both sorption and permeation measurements.

**Gases.** The gases and vapors used in the permeation and sorption experiments had a purity of at least 99.5%. N<sub>2</sub> and H<sub>2</sub> were obtained from National Specialty Gases (Durham, NC) while C<sub>3</sub>H<sub>8</sub> and C<sub>3</sub>F<sub>8</sub> were purchased from Machine Welding (Raleigh, NC). All gases were used as received.

**Characterization.** Sorption. Solubility coefficients were determined using a high-pressure barometric apparatus.<sup>33</sup> Initially, a polymer film was placed in the sample chamber and exposed to vacuum overnight to remove the air gases. A known amount of penetrant gas was introduced into the chamber, and the pressure was allowed to equilibrate. Once

the chamber pressure was constant, the amount of gas sorbed by the polymer was determined by mass balance. Then additional penetrant was introduced, and the procedure was repeated. In this incremental manner, penetrant uptake was determined as a function of pressure. The maximum pressure was 3–25 atm depending on the penetrant. Sorption equilibrium for all gases was reached within, at most, a few hours. After measuring each isotherm, the polymer samples were degassed under vacuum overnight. The system temperature was controlled to  $\pm 0.1$  °C using a constant temperature water bath. For PDMS, the sorption experiments were performed in the following order: N<sub>2</sub>, H<sub>2</sub>, C<sub>3</sub>H<sub>8</sub>, and C<sub>3</sub>F<sub>8</sub>. For the last two gases, solubility was measured at different temperatures in the order of increasing temperature, i.e., 25, 35, 45, and then 55 °C. The orders of gases and temperatures for PTMSP were also the same, except that sorption of H<sub>2</sub> in PTMSP was not measured.

**Permeation.** Pure gas permeability coefficients were determined using a constant pressure/variable volume apparatus.<sup>34</sup> The membrane area was 13.8 cm<sup>2</sup>. The upstream pressure was varied from 2 to 17.4 atm for N<sub>2</sub> and H<sub>2</sub> while for C<sub>3</sub>H<sub>8</sub> and C<sub>3</sub>F<sub>8</sub> it was kept constant at 2.36 atm. In all cases, the downstream pressure was atmospheric. Gas flow rates were measured with a soap-film bubble flowmeter. The system temperature was controlled to  $\pm 0.5$  °C using a DYNA-SENSE temperature control system. Prior to each experiment, the upstream and downstream sides of the permeation cell were purged with penetrant gas. Permeability coefficients of the gases and vapors were determined in the following order: N<sub>2</sub>, H<sub>2</sub>, C<sub>3</sub>H<sub>8</sub>, and C<sub>3</sub>F<sub>8</sub> where, for each gas, measurements at different temperatures were made in the order of increasing temperature. For PTMSP, to minimize conditioning effects, a fresh film was used for each gas. The variation in nitrogen permeability from film to film, at all temperatures measured, was less than 10%. When pseudo-steady-state conditions were attained, the following expression was used to evaluate permeability:

$$P = \frac{22414}{A} \frac{l}{p_2 - p_1} \frac{p_1}{RT} \frac{dV}{dt} \quad (16)$$

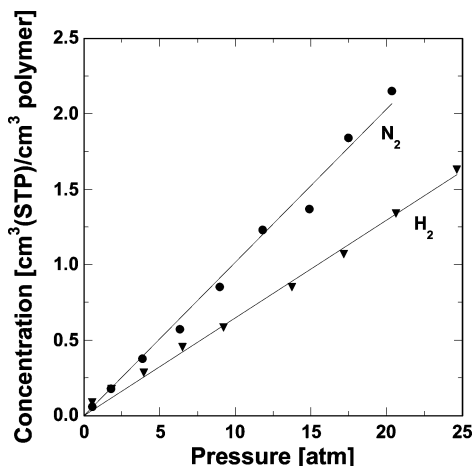
where  $A$  is the membrane area (cm<sup>2</sup>),  $l$  is the membrane thickness (cm),  $p_2$  is the upstream pressure (atm),  $p_1$  is the downstream pressure (atmospheric pressure in this case),  $R$  is the universal gas constant (6236.56 cm<sup>3</sup> cmHg/(mol K)),  $T$  is the absolute temperature (K), and  $dV/dt$  is the volumetric displacement rate of the soap film in the bubble flowmeter (cm<sup>3</sup>/s).

## Results and Discussion

**Solubility.** PDMS. Sorption isotherms for nitrogen and hydrogen in PDMS at 35 °C are presented in Figure 1. The isotherms obey Henry's law, and our experimental data are in good agreement with previously published data for nitrogen sorption in PDMS at 35 °C.<sup>35</sup> From Figure 1, the ratio of nitrogen to hydrogen solubility is approximately 1.6. The value of this ratio in a wide variety of liquids lies between 1.2 and 2.2.<sup>36</sup> For example, the N<sub>2</sub>/H<sub>2</sub> solubility ratio is 1.4 in carbon disulfide, around 1.7 in alcohols, and in the range 1.9–2.2 in hydrocarbon liquids at 25 °C and 1 atm.<sup>36</sup> Thus, the N<sub>2</sub>/H<sub>2</sub> solubility ratio in PDMS lies in the same range as that in liquids, and this is one simple method for ensuring that the data are reasonable, since this ratio is expected to be comparable among rubbery polymers and liquids.

Sorption isotherms for propane and perfluoropropane in PDMS at 25, 35, 45, and 55 °C are presented in parts a and b of Figure 2, respectively. Perfluoropropane solubility in PDMS is enormously lower than that of propane. For example, at 35 °C and 3 atm the sorbed





**Figure 1.** N<sub>2</sub> and H<sub>2</sub> sorption in PDMS at 35 °C.

concentrations of C<sub>3</sub>H<sub>8</sub> and C<sub>3</sub>F<sub>8</sub> are 23 and 2.3 cm<sup>3</sup> (STP)/(cm<sup>3</sup> polymer), respectively, a difference of 1 order of magnitude. For both penetrants, solubility decreases with increasing temperature at a given pressure, indicating that the sorption process is exothermic.

Propane sorption isotherms are convex to the pressure axis, which is consistent with the behavior of highly sorbing penetrants in rubbery polymers.<sup>27</sup> The curvature of the isotherms decreases with increasing temperature, suggesting a weaker dependence of solubility on pressure at higher temperatures. This is consistent with the findings of Shah et al.,<sup>37</sup> who observed a decrease in the pressure dependence of propane solubility in PDMS as temperature increased. They obtained an infinite dilution solubility of 6.45 cm<sup>3</sup> (STP)/(cm<sup>3</sup> polymer atm) at 35 °C, which is in excellent agreement with our value of 6.5 ± 0.06 cm<sup>3</sup> (STP)/(cm<sup>3</sup> polymer atm). If the propane sorption isotherms are plotted as a function of activity (i.e.,  $p/p_{\text{sat}}$ ) instead of pressure (cf. Figure 2c), the four isotherms collapse to a single curve. This result suggests that the change in the curvature of the isotherms with temperature is a result of exploring a smaller activity range at higher temperatures, since the maximum pressure in these experiments is almost the same but the value of  $p_{\text{sat}}$  increases substantially with temperature.

Perfluoropropane sorption isotherms are linear (cf. Figure 2b). When the amount of perfluoropropane sorbed in PDMS is plotted against penetrant activity (cf. Figure 2d), which should account for variability in C<sub>3</sub>F<sub>8</sub> condensability with temperature, the C<sub>3</sub>F<sub>8</sub> sorbed concentration increases with temperature. This behavior is qualitatively unlike that of C<sub>3</sub>H<sub>8</sub> and suggests less favorable interactions between C<sub>3</sub>F<sub>8</sub> and the PDMS matrix than between C<sub>3</sub>H<sub>8</sub> and PDMS.

In the absence of specific interactions between penetrant molecules and the polymer matrix, solubility coefficients usually scale with measures of penetrant condensability such as critical temperature,  $T_c$ .<sup>27</sup> Such relationships often utilize solubility coefficients in the limit of zero pressure (called infinite dilution solubility,  $S^\infty$ ) to compare solubilities of penetrants on a consistent basis:

$$S^\infty = \lim_{p \rightarrow 0} S = \lim_{p \rightarrow 0} \frac{C}{p} \quad (17)$$

Suwandi and Stern<sup>38</sup> observed a linear correlation of the logarithm of infinite dilution solubility,  $S^\infty$ , with

$(T_c/T)^2$  for a large number of penetrants in PDMS. This result has been reproduced in Figure 3 (the numerical data have been tabulated previously<sup>10</sup>) along with the  $S^\infty$  values of propane and perfluoropropane determined in this study (filled symbols). The data for propane obey this correlation, but perfluoropropane solubility coefficients fall well below the trendline, consistent with our previous isothermal study of hydrocarbon and fluorocarbon solubility in PDMS.<sup>10</sup> This result suggests that the fluorinated penetrant experiences less favorable interactions with PDMS than its hydrocarbon analogue.

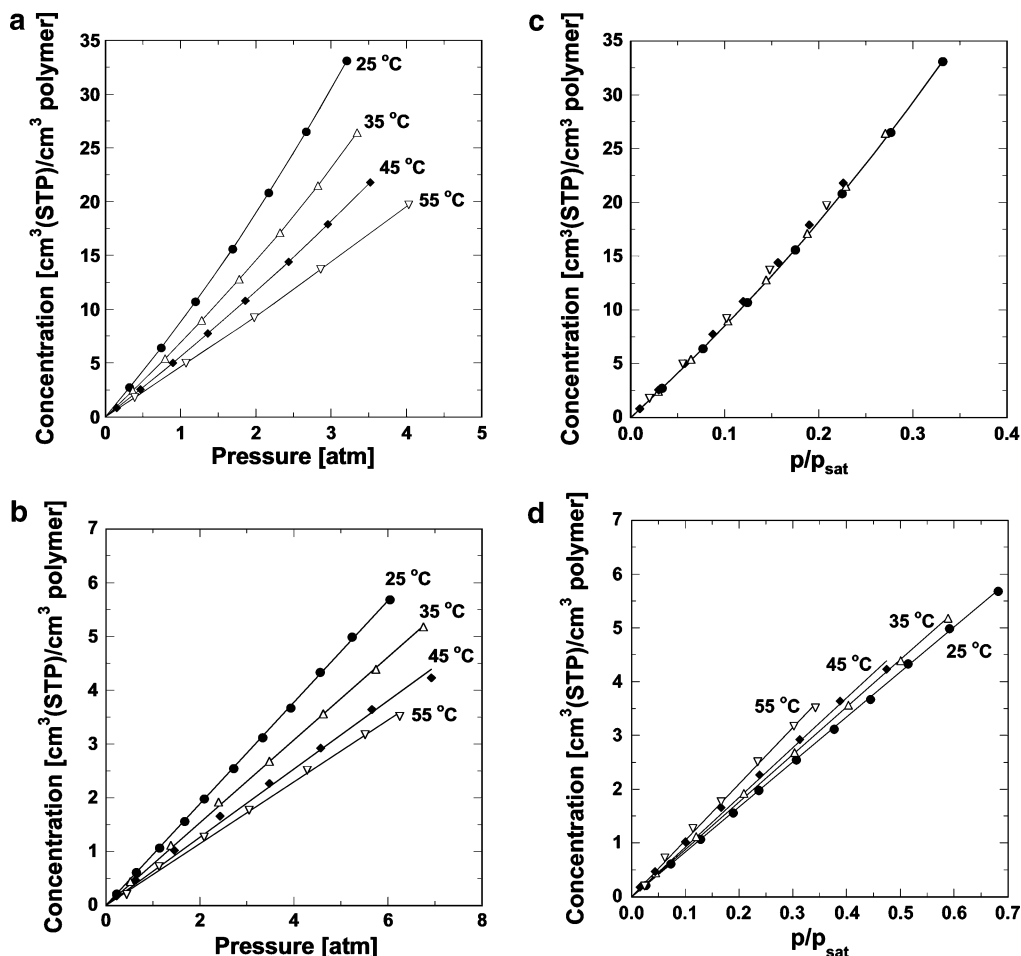
From sorption data such as those presented in Figure 2a,b, the effect of temperature on solubility can be determined. Solubility values from these figures, calculated either at constant penetrant pressure or at a constant penetrant concentration in the polymer, can be fitted to the van't Hoff equation (eq 14) to obtain enthalpies of penetrant sorption. Thus, two enthalpies of sorption can be calculated viz. the enthalpy of sorption at constant pressure (isobaric),  $\Delta H_s^P$ , or the enthalpy of sorption at constant concentration (isosteric),  $\Delta H_s^C$ . The relationship between the isobaric and isosteric enthalpies of sorption is derived in the Appendix. In this document,  $\Delta H_s$  will be used in equations which apply to both isosteric and isobaric enthalpies of sorption, while the superscripts "C" and "P" will be used to denote the particular types of enthalpies when it is necessary to make this distinction.

Isosteric enthalpies of sorption were calculated for C<sub>3</sub>H<sub>8</sub> and C<sub>3</sub>F<sub>8</sub> in PDMS from the data in Figure 2a,b. Since sorption is typically viewed as a two-step process involving penetrant condensation from a gaslike density to a liquidlike density followed by mixing condensed penetrant molecules with polymer segments, the enthalpy of sorption can be viewed as a sum of the enthalpy changes for these two steps:<sup>39</sup>

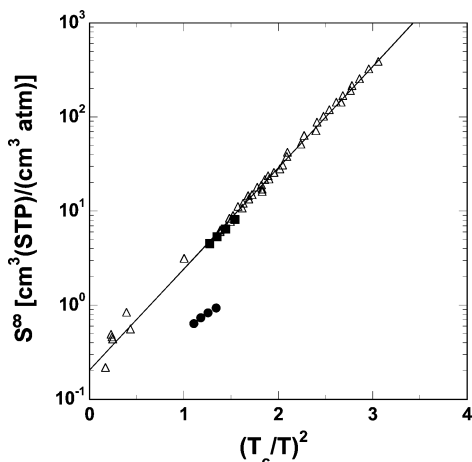
$$\Delta H_s = \Delta H_{\text{cond}} + \Delta H_m \quad (18)$$

where  $\Delta H_{\text{cond}}$  and  $\Delta H_m$  are the enthalpy changes associated with penetrant condensation and mixing, respectively.<sup>39</sup> To estimate  $\Delta H_m$ , a value of  $\Delta H_{\text{cond}}$  must be supplied. However,  $\Delta H_{\text{cond}}$  varies somewhat over the temperature range of study. For example,  $\Delta H_{\text{cond}}$  for C<sub>3</sub>H<sub>8</sub> and C<sub>3</sub>F<sub>8</sub> changes by 2.7 and 4.5 kJ/mol, respectively, over the temperature range investigated. For the present calculation,  $\Delta H_{\text{cond}}$  values have been taken at 40 °C, which is the midpoint of the experimental temperature range. The  $\Delta H_{\text{cond}}$  values are −13.6 and −12.7 kJ/mol for C<sub>3</sub>H<sub>8</sub> and C<sub>3</sub>F<sub>8</sub>, respectively, at this temperature.<sup>40</sup> These values were used in eq 18 to calculate the isosteric enthalpy of mixing in PDMS, and the results are presented in Figure 4a. Basing the calculation of  $\Delta H_m$  on the  $\Delta H_{\text{cond}}$  value at 40 °C is an arbitrary choice, and the values of  $\Delta H_m$  in Figure 4a will vary somewhat depending on the value of  $\Delta H_{\text{cond}}$  used. However, there is no qualitative change in the order of the data presented in Figure 4a or its dependence on concentration if other reasonable reference temperatures are used for determining  $\Delta H_{\text{cond}}$ , so the discussion and conclusions below are not affected by this choice.

As indicated in Figure 4a, the enthalpy of mixing perfluoropropane with PDMS segments is much greater than that of propane. For example, at infinite dilution the enthalpies of mixing for C<sub>3</sub>F<sub>8</sub> and C<sub>3</sub>H<sub>8</sub> are 2.5 and −2.8 kJ/mol, respectively. In both cases, the enthalpy



**Figure 2.** (a)  $C_3H_8$  and (b)  $C_3F_8$  sorption in PDMS as a function of temperature. (c)  $C_3H_8$  and (d)  $C_3F_8$  sorption in PDMS as a function of penetrant activity ( $p/p_{\text{sat}}$ ) at four temperatures: (●) 25, (△) 35, (◆) 45, and (▽) 55 °C.  $p_{\text{sat}}$  values are from the correlations in Appendix A of Reid et al.<sup>52</sup>



**Figure 3.** Correlation of infinite dilution solubility,  $S^\infty$ , in PDMS with reduced critical temperature: (■) propane data of this study; (●) perfluoropropane data of this study; (△) data of Suwandi and Stern,<sup>38</sup> Barrer et al.,<sup>53</sup> and Robb.<sup>54</sup> The correlation line is  $S^\infty [\text{cm}^3 (\text{STP})/(\text{cm}^3 \text{ atm})] = 0.0245 \times 10^{1.075(T_c/T)^2}$ .

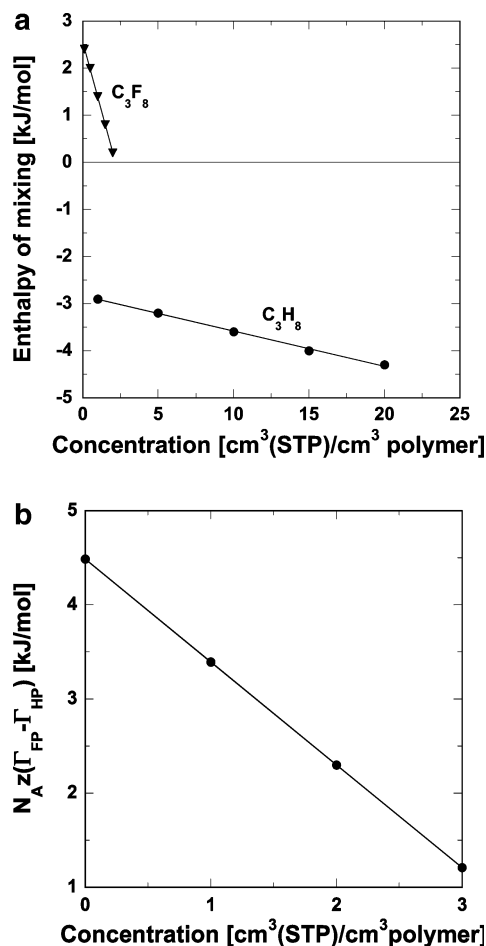
of mixing decreases with concentration, implying that the process of mixing becomes more favorable in the presence of greater amounts of penetrant. Typically, if the polymer matrix and the penetrant molecules are chemically dissimilar and do not have specific interactions with each other, interactions among penetrant molecules are more favorable than those between pen-

etrant molecules and polymer chains. At low penetrant concentrations, mixing these penetrant molecules with the polymer is a less favorable process than at higher concentrations. As penetrant concentration increases, the environment into which the penetrant is dissolving becomes more like that of the penetrant, and the mixing process becomes more favorable. The enthalpy of mixing of propane depends much less on penetrant concentration than that of perfluoropropane. This result is reasonable because, from a structural viewpoint, propane and PDMS are much more similar than perfluoropropane and PDMS.

The difference in interaction energy of PDMS with perfluoropropane and propane can be estimated from polymer–polymer, penetrant–penetrant, and polymer–penetrant interaction energies. Based on the regular solution and Flory–Huggins theories<sup>10</sup>

$$z(\Gamma_{\text{FP}} - \Gamma_{\text{HP}}) = \frac{RT}{N_A} \left[ (\chi_F - \chi_H) - \frac{1}{RT} (\Delta H_{\text{condH}} - \Delta H_{\text{condF}}) \right] \quad (19)$$

where  $\Gamma_{\text{FP}}$  is the potential energy required to separate a perfluoropropane molecule and a PDMS segment to infinite distance,  $\Gamma_{\text{HP}}$  is the potential energy required to separate a propane molecule and a PDMS segment to infinite distance,  $z$  is the coordination number,  $N_A$  is Avogadro's number,  $\chi_F$  and  $\chi_H$  are the Flory–Huggins interaction parameters for perfluoropropane and pro-



**Figure 4.** (a) Isosteric enthalpy of mixing of C<sub>3</sub>H<sub>8</sub> and C<sub>3</sub>F<sub>8</sub> in PDMS. (b) Difference in potential energies associated with insertion of C<sub>3</sub>F<sub>8</sub> and C<sub>3</sub>H<sub>8</sub> in PDMS.

pane, respectively, and  $\Delta H_{\text{condF}}$  and  $\Delta H_{\text{condH}}$  are the molar enthalpies of condensation of perfluoropropane and propane, respectively. The product  $z(\Gamma_{FP} - \Gamma_{HP})$  is the difference in potential energy associated with inserting a C<sub>3</sub>F<sub>8</sub> molecule and a C<sub>3</sub>H<sub>8</sub> molecule in PDMS. Also<sup>41</sup>

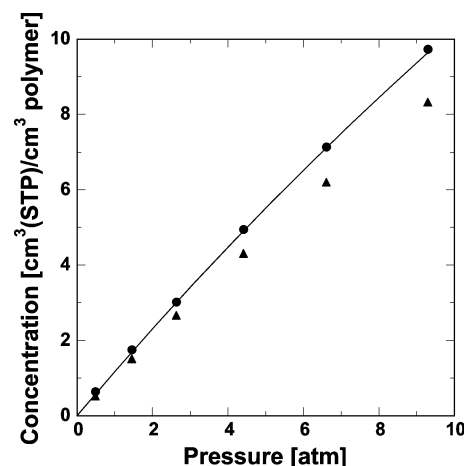
$$\Delta H_m = RT\chi(1 - \phi_2) \quad (20)$$

where  $\Delta H_m$  is the enthalpy of mixing and  $\phi_2$  is the penetrant volume fraction (which is  $\ll 1$  in this study). Combining eqs 19 and 20 yields

$$\begin{aligned} N_A z(\Gamma_{FP} - \Gamma_{HP}) &= (\Delta H_{mF} - \Delta H_{mH}) - (\Delta H_{\text{condH}} - \Delta H_{\text{condF}}) \\ &= \Delta H_{sF}^C - \Delta H_{sH}^C \end{aligned} \quad (21)$$

where  $\Delta H_{sF}^C$  and  $\Delta H_{sH}^C$  are the isosteric enthalpies of mixing the fluorocarbon and hydrocarbon penetrants, respectively, with the polymer segments. The left-hand side of eq 21 is the difference in energy associated with inserting a mole of C<sub>3</sub>F<sub>8</sub> molecules and a mole of C<sub>3</sub>H<sub>8</sub> molecules in PDMS.

Figure 4b presents the calculated difference in potential energy associated with inserting C<sub>3</sub>F<sub>8</sub> molecules and C<sub>3</sub>H<sub>8</sub> molecules in PDMS. In the limit of infinite dilution, this difference is 4.5 kJ/mol while at the highest concentration considered, 3 cm<sup>3</sup> (STP)/(cm<sup>3</sup> polymer), it is 1.2 kJ/mol. The result at infinite dilution



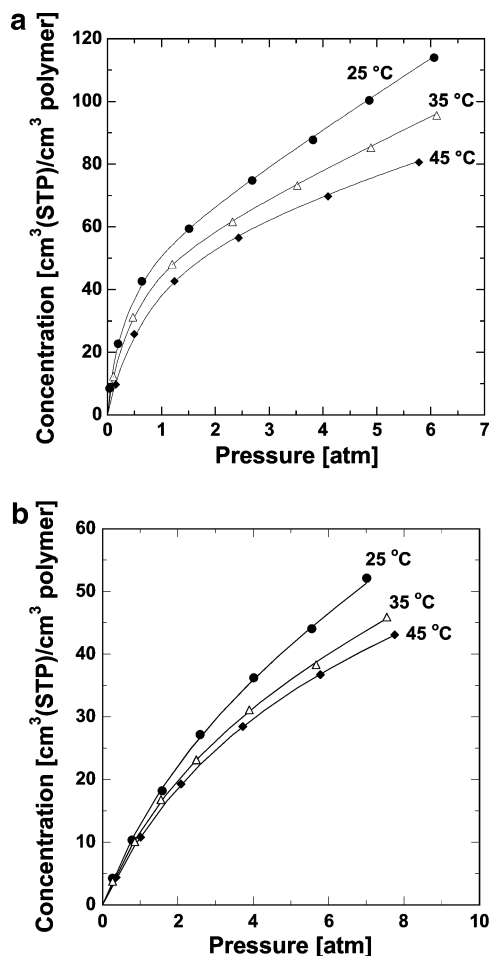
**Figure 5.** N<sub>2</sub> sorption in PTMSP at 35 °C. Data of Ichiraku et al.<sup>42</sup> (▲) are provided for comparison.

is in excellent agreement with that previously estimated from isothermal sorption data (4.6 kJ/mol at infinite dilution).<sup>10</sup>

**PTMSP.** A nitrogen sorption isotherm in PTMSP at 35 °C is presented in Figure 5 along with previously published data for comparison.<sup>42</sup> Our data are in good agreement with the literature data. Sorption isotherms for propane and perfluoropropane in PTMSP at 25, 35, and 45 °C are presented in parts a and b of Figure 6, respectively. The isotherms are concave to the pressure axis, which is typical for gas sorption in glassy polymers.<sup>27</sup> There is a substantial difference in the solubilities of the hydrocarbon and fluorocarbon analogues, with propane being more soluble in PTMSP at all temperatures and pressures tested. As noted previously,<sup>9</sup> this difference in hydrocarbon and perfluorocarbon solubilities is smaller in high free volume PTMSP than in liquidlike PDMS. For example, at 35 °C and 3 atm, the sorbed concentrations of C<sub>3</sub>H<sub>8</sub> and C<sub>3</sub>F<sub>8</sub> in PTMSP are 68 and 26 cm<sup>3</sup> (STP)/(cm<sup>3</sup> polymer), respectively, which is significantly smaller than the order of magnitude difference observed in PDMS.

Isosteric enthalpies of sorption were calculated for C<sub>3</sub>H<sub>8</sub> and C<sub>3</sub>F<sub>8</sub> in PTMSP from the data in Figure 6a,b. The enthalpies of condensation for C<sub>3</sub>H<sub>8</sub> (−14 kJ/mol) and C<sub>3</sub>F<sub>8</sub> (−13.4 kJ/mol) were taken at 35 °C, the midpoint of the experimental temperature range.<sup>40</sup> The isosteric enthalpies of mixing in PTMSP were then calculated from eq 18 and are presented in Figure 7a as a function of penetrant concentration. Similar to PDMS, the enthalpy of mixing of C<sub>3</sub>F<sub>8</sub> in PTMSP is generally more positive (i.e., endothermic) than that of C<sub>3</sub>H<sub>8</sub> in PTMSP. However, unlike PDMS, the enthalpies of mixing of the two penetrants show opposite trends in PTMSP; the enthalpy of mixing C<sub>3</sub>F<sub>8</sub> with PTMSP decreases with concentration while the enthalpy of mixing C<sub>3</sub>H<sub>8</sub> with PTMSP increases with concentration.

The trend of the enthalpy of mixing of propane can be rationalized by considering the additional mode of sorption available in glassy polymers, i.e., the Langmuir microvoids. A glassy polymer such as PTMSP contains nonequilibrium microvoids (so-called Langmuir sites) dispersed throughout the equilibrium matrix. These microvoids represent sorption sites that, from an energetic perspective, are easily accessible to penetrant

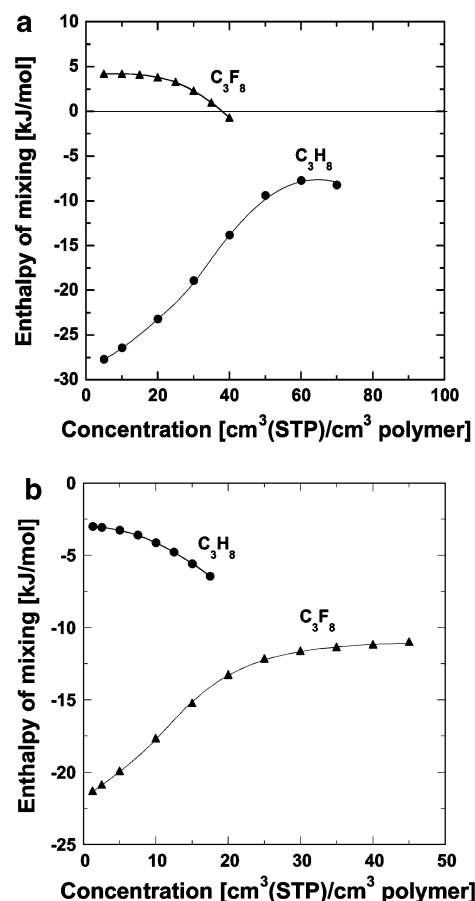


**Figure 6.** (a)  $C_3H_8$  and (b)  $C_3F_8$  sorption in PTMSP as a function of temperature.

molecules. At low penetrant concentrations, the microvoids are relatively unoccupied and can accommodate penetrant molecules with little or no distortion (i.e., swelling) of the polymer matrix. As penetrant concentration in the polymer increases, these sites become progressively more saturated, resulting in a larger fraction of the sorption occurring in the densified regions of the polymer (i.e., the so-called Henry's law region). Penetrant sorption in the Henry's law region is energetically more expensive than sorption in a Langmuir microvoid, since this process involves the creation of a gap large enough to accommodate the penetrant (i.e., the polymer swells<sup>43</sup>). Thus, the enthalpy of mixing increases with concentration. At high concentrations, it reaches a limiting value determined by the enthalpy of sorption of the penetrant in the densified matrix. A similar trend has been reported for  $CO_2$  sorption in a high barrier, glassy polymer, poly(ethylene terephthalate).<sup>44</sup>

The concentration dependence of the enthalpy of mixing of perfluoropropane in PTMSP is similar to that in PDMS. This result suggests that the less favorable interactions between the fluorinated penetrant and the hydrocarbon matrix are more important than the dual-mode effects in determining sorption energetics.

An interesting example of mixing behavior that follows the same logic is displayed in Figure 7b for isosteric enthalpies of mixing  $C_3H_8$  and  $C_3F_8$  with a fluorinated copolymer of 87 mol % 2,2-bis(trifluoromethyl)-4,5-difluoro-1,3-dioxole and 13 mol % tetra-



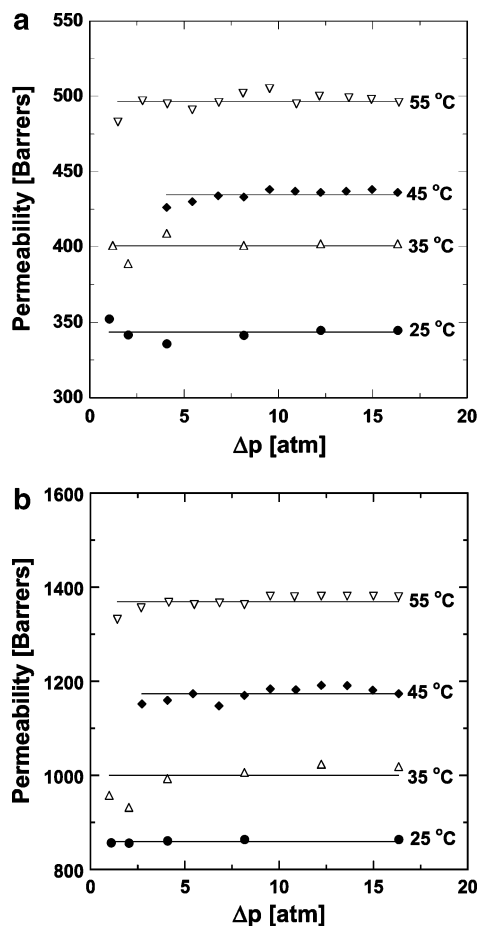
**Figure 7.** Isosteric enthalpies of mixing of  $C_3H_8$  and  $C_3F_8$  in (a) PTMSP and (b) Teflon AF2400.<sup>8</sup>

fluoroethylene, called Teflon AF2400.<sup>8</sup> In this case, the penetrants encounter a fluorinated environment, and the concentration dependence of the enthalpies of mixing as well as their relative magnitudes reflect this fact. Propane has a less exothermic enthalpy of mixing, and  $\Delta H_m$  decreases with concentration, both of which signify less favorable polymer–penetrant interactions. On the other hand, the concentration dependence of perfluoropropane's enthalpy of mixing suggests that Langmuir microvoid filling is the dominant factor affecting the energetics of the dissolution process.

**Permeability. PDMS.** The permeabilities of PDMS to nitrogen and hydrogen in PDMS as a function of the pressure difference across the polymer film at 25, 35, 45, and 55 °C are shown in parts a and b of Figure 8, respectively. The permeabilities of both penetrants increase with increasing temperature at a given pressure, indicating a positive activation energy of permeation. This result is typical for the permeation of supercritical gases in PDMS. For example, Bixler and Sweeting<sup>45</sup> reported an  $E_P$  value for nitrogen in PDMS of 10.9 kJ/mol, which is similar to the value of 9.3 kJ/mol obtained from the data in Figure 8a.

Parts a and b of Figure 9 present the effect of temperature on the permeability of PDMS to propane and perfluoropropane, respectively, at an upstream pressure of 2.36 atm. The data can be fit to the Arrhenius equation (eq 12), and the activation energy of permeation,  $E_P$ , can be calculated. Since the permeability values are at a constant upstream pressure, the  $E_P$  value thus calculated is an isobaric (i.e., constant pressure) activation energy of permeation. In a similar





**Figure 8.** (a)  $N_2$  and (b)  $H_2$  permeation in PDMS as a function of temperature and pressure difference across the membrane. The downstream pressure is 1 atm.

**Table 1. Activation Energies of Permeation and Diffusion, and Enthalpy of Sorption at 2.36 atm (i.e., Isobaric) for  $C_3H_8$  and  $C_3F_8$  in PDMS and PTMSP<sup>a</sup>**

polymer	penetrant	$E_P$ (kJ/mol)	$\Delta H_s^P$ (kJ/mol)	$E_D$ (kJ/mol)
PDMS	$C_3H_8$	$-13 \pm 1.2$	$-20 \pm 0.2$	$7 \pm 1.2$
	$C_3F_8$	$3 \pm 3$	$-13 \pm 1.8$	$16 \pm 3$
PTMSP	$C_3H_8$	$-8 \pm 1.2$	$-9 \pm 0.7$	$1 \pm 1.4$
	$C_3F_8$	$7 \pm 4.0$	$-7 \pm 2.0$	$14 \pm 4.5$

<sup>a</sup> Note:  $E_P$ ,  $\Delta H_s$ , and  $E_D$  values have been calculated using eqs 12, 14, and 15, respectively.

fashion, isobaric enthalpies of sorption can be calculated for the two penetrants at a pressure of 2.36 atm, from eq 14 and the data in Figure 2a,b. The calculated values are presented in Table 1. The errors in the table have been calculated by the method of propagation of errors.<sup>46</sup> These  $E_P$  and  $\Delta H_s^P$  values have been used to calculate  $E_D$  values according to eq 15. As indicated earlier, the use of eq 15 is subject to restrictions on the concentration or pressure dependence of the permeability, solubility, and diffusion coefficients. However, since the permeability data have been measured at a single pressure, it is not possible to evaluate the pressure (or concentration) dependence of the permeability coefficients. Therefore, the activation energy of diffusion has been calculated by assuming the validity of eq 15, as is practically always done in the literature.

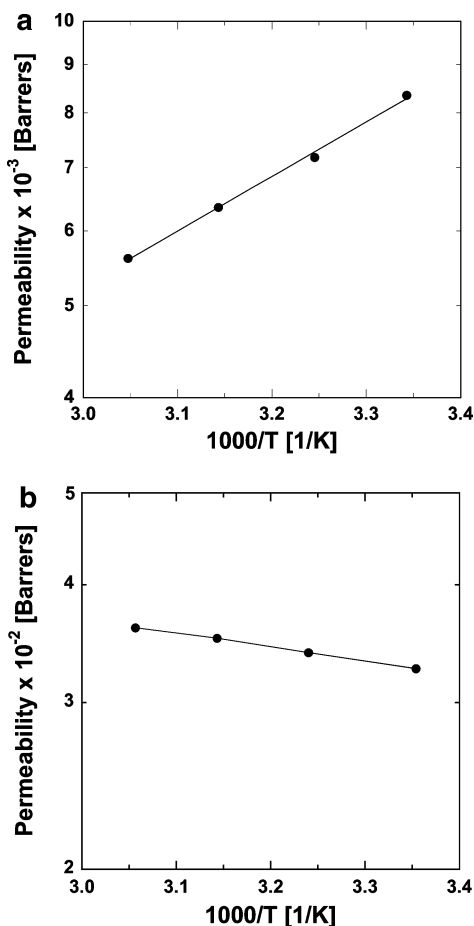
Usually, the permeability of relatively condensable gases in PDMS decreases with increasing temperature (i.e.,  $E_P$  is negative).<sup>47</sup> This is because, in weakly size sieving PDMS, the solubility of condensable gases

decreases with increasing temperature more rapidly than diffusion coefficients increase (i.e.,  $|\Delta H_s| > E_D$ ). The propane data in PDMS are consistent with this trend. However, perfluoropropane displays markedly different behavior. Its permeability coefficient increases slightly with temperature, indicating a positive  $E_P$  value. From Table 1, the different effect of temperature on the permeation behavior of the  $C_3$  analogues is related to differences in both the dissolution and diffusion of these penetrants. Larger penetrants require more energy to execute a diffusive jump in a polymer matrix than smaller penetrants.<sup>48</sup> In this regard, perfluoropropane is substantially larger than its hydrocarbon analogue. The critical volumes of  $C_3F_8$  and  $C_3H_8$  are 300 and 203  $cm^3/mol$ , respectively. Consequently, the activation energy of diffusion is much larger for  $C_3F_8$  (16 kJ/mol) than for  $C_3H_8$  (7 kJ/mol). Additionally, a difference in polymer–penetrant interactions contributes to the difference in the temperature dependence of permeability of  $C_3F_8$  and  $C_3H_8$ . From Table 1, the magnitude of the isobaric enthalpy of sorption is lower for  $C_3F_8$  than for  $C_3H_8$ . The enthalpies of condensation for  $C_3F_8$  and  $C_3H_8$  are similar ( $-12.7$  and  $-13.6$  kJ/mol, respectively, at 40 °C<sup>40</sup>), so the mixing of propane with PDMS is more favorable than mixing perfluoropropane with PDMS. The combination of less favorable mixing and hindered diffusion causes  $C_3F_8$  to have a substantially larger activation energy of permeation than  $C_3H_8$  in PDMS.

**PTMSP.** Figure 10a,b presents  $N_2$  and  $H_2$  permeability in PTMSP as a function of temperature and pressure. For these penetrants, permeability coefficients decrease with increasing temperature, which is opposite to the behavior in PDMS. This temperature dependence of permeability is consistent with the high free volume and poor size-sieving ability of PTMSP.

Parts a and b of Figure 11 present Arrhenius plots of  $C_3H_8$  and  $C_3F_8$  permeability coefficients, respectively, in PTMSP at 2.36 atm. For  $C_3H_8$ , permeability increases as temperature decreases, which is qualitatively similar to the trend in PDMS. On the other hand,  $C_3F_8$  permeability decreases as temperature decreases, and this is the first report of permeability coefficients decreasing with decreasing temperature in PTMSP. Isobaric activation energies of permeation for these penetrants are tabulated in Table 1 along with enthalpies of sorption and calculated activation energies of diffusion. Similar to PDMS,  $E_P$  for  $C_3F_8$  in PTMSP is positive and opposite in sign to that of propane. This result reflects the larger size of perfluoropropane (larger  $E_D$ ) and its less favorable interactions with PTMSP (larger  $\Delta H_m$  and, therefore, more positive  $\Delta H_s$  relative to propane). The relative contributions of these two effects to the difference in  $E_P$  between  $C_3H_8$  and  $C_3F_8$  in PTMSP are shown in Table 1. For PTMSP, the difference in  $C_3H_8$  and  $C_3F_8$  activation energies of permeation (15 kJ/mol) is mostly due to the large difference in  $E_D$  values (13 kJ/mol) with a small contribution from the  $\Delta H_s$  difference (2 kJ/mol). Thus, for PTMSP, it is primarily the difference in penetrant size and the associated effect on the diffusion process that causes the dramatic difference in  $C_3$  analogue transport properties. In contrast, in PDMS the difference in  $E_P$  values between  $C_3H_8$  and  $C_3F_8$  (16 kJ/mol) is due to nearly equal contributions from the difference in  $E_D$  (9 kJ/mol) and  $\Delta H_s$  (7 kJ/mol). This result suggests that hydrocarbon–fluorocarbon interac-



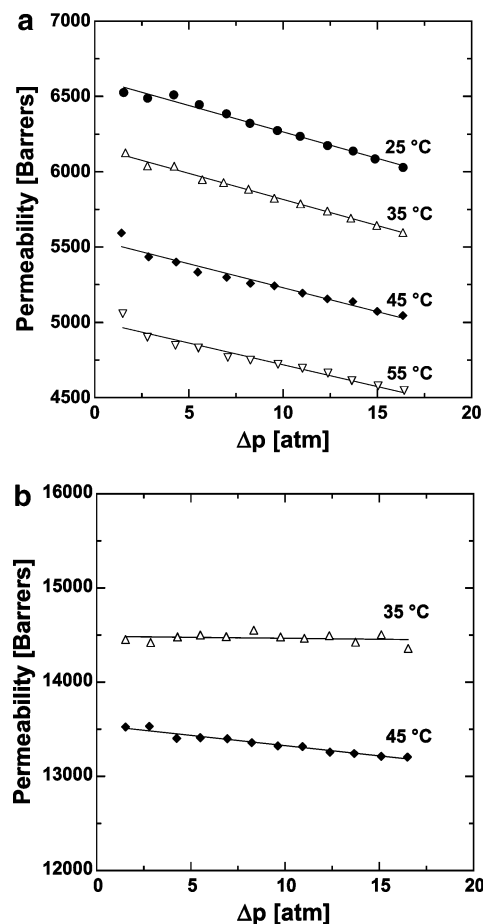


**Figure 9.** Effect of temperature on (a) C<sub>3</sub>H<sub>8</sub> and (b) C<sub>3</sub>F<sub>8</sub> permeation in PDMS at 2.36 atm upstream pressure and 1 atm downstream pressure.

tions have a stronger effect on penetrant transport properties in liquidlike PDMS than in high free volume, glassy PTMSP.

The unusual nature of C<sub>3</sub>F<sub>8</sub>'s positive activation energy of permeation in PTMSP is highlighted in Figure 12, which presents  $E_P$  values as a function of penetrant critical temperature for this polymer. The data of Masuda et al.<sup>49</sup> are included in Figure 12 for comparison. Excluding our data for C<sub>3</sub>F<sub>8</sub>,  $E_P$  values in PTMSP are negative for all penetrants and decrease (i.e., increase in absolute value) with increasing penetrant critical temperature. This behavior is consistent with a solubility selective polymer where  $E_P$  is significantly influenced by  $\Delta H_s$  (because  $|\Delta H_s| > E_D$ ) and where the enthalpy of sorption decreases with increasing penetrant condensability ( $T_c$ ). The activation energy of permeation for C<sub>3</sub>F<sub>8</sub> deviates substantially from the empirical trendline through the other penetrants. In fact, based on C<sub>3</sub>F<sub>8</sub>'s critical temperature and the best fit line through the rest of the data, its expected  $E_P$  value would be  $-9.8$  kJ/mol as compared to the measured value of  $7$  kJ/mol. As mentioned previously, this difference appears to be primarily due to C<sub>3</sub>F<sub>8</sub>'s large size and its effect on the diffusion process.

It has been suggested<sup>5,50</sup> that PTMSP contains a network of quasi-permanent, interconnected free volume elements spanning the polymer through which the majority of penetrant transport occurs. Transport in these interconnected free volume elements may be similar to that in zeolites, where the critical penetrant diameter, or the smallest size window through which a



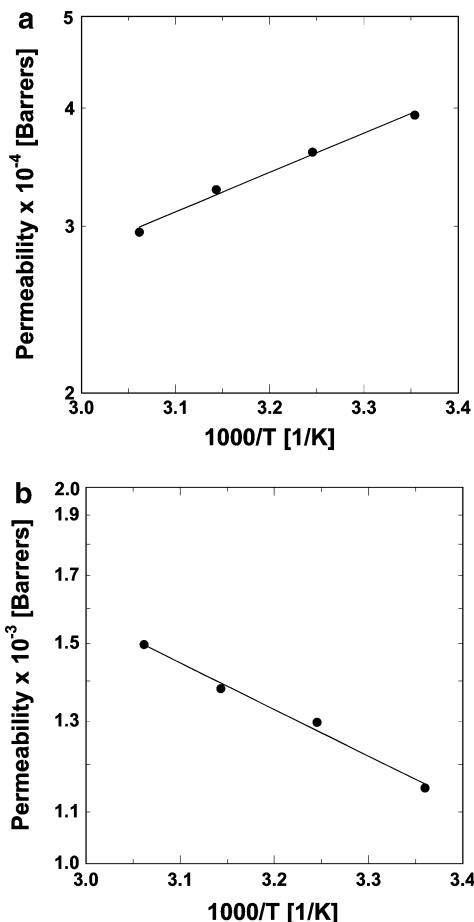
**Figure 10.** (a) N<sub>2</sub> and (b) H<sub>2</sub> permeation in PTMSP as a function of temperature and pressure difference across the membrane. The downstream pressure is 1 atm.

given molecule can fit, governs the transport. Transport through such interconnected free volume elements should be an energetically inexpensive process compared to transport through the densified polymer matrix. The kinetic diameter of C<sub>3</sub>F<sub>8</sub> (5.4 Å) is much larger than that of C<sub>3</sub>H<sub>8</sub> (4.3 Å).<sup>51</sup> Thus, it is possible that C<sub>3</sub>F<sub>8</sub> is larger than the critical free volume element diameter for transport in the interconnected free volume elements, which may restrict its access to this energetically inexpensive mode of transport.

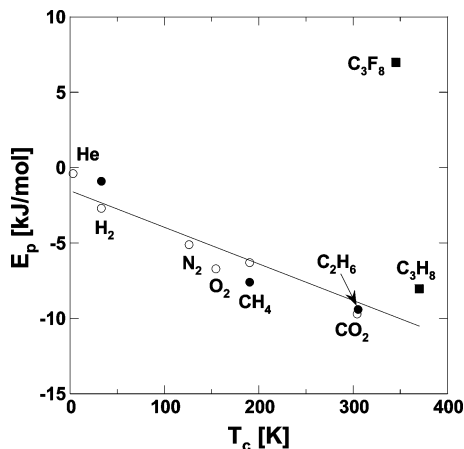
## Conclusions

C<sub>3</sub>F<sub>8</sub> solubility is lower than that of its hydrocarbon analogue, C<sub>3</sub>H<sub>8</sub>, in both PDMS and PTMSP due to less favorable polymer–fluorocarbon interactions as indicated by a higher  $\Delta H_m$  for the fluorocarbon as well as a decrease in  $\Delta H_m$  with increasing penetrant concentration. Unlike PDMS,  $\Delta H_m$  of C<sub>3</sub>H<sub>8</sub> in PTMSP increases with penetrant concentration, which is due to the different energy requirements for dissolution in the two types of sorption sites present in a glassy polymer viz. the Langmuir sites and the Henry's law sites.

The activation energy of permeation for C<sub>3</sub>F<sub>8</sub> is positive and opposite in sign to that of C<sub>3</sub>H<sub>8</sub> in both polymers. This is the first report of a penetrant having a positive activation energy of permeation in PTMSP. In PDMS, the difference in  $E_P$  values between propane and perfluoropropane is due to the difference in penetrant sizes as well as the difference in polymer–penetrant interactions. For PTMSP, however, it is primarily the larger size of C<sub>3</sub>F<sub>8</sub> and its associated effect



**Figure 11.** Effect of temperature on (a)  $C_3H_8$  and (b)  $C_3F_8$  permeation in PTMSP at 2.36 atm upstream pressure and 1 atm downstream pressure.



**Figure 12.** Activation energy of permeation of various penetrants in PTMSP: (●) data from Masuda et al.;<sup>49</sup> (○) unpublished data of T. C. Merkel and Z. He at Membrane Technology and Research, Inc. (Menlo Park, CA); (■) data of this study. The straight-line in the figure is the least-squares fit to the data for all the penetrants except  $C_3F_8$  and is given by  $E_p$  [kJ/mol] =  $-1.52 - 0.024 \times T_c$  [K].

on diffusion that is responsible for the difference in  $E_p$  values. Thus, polymer–penetrant interactions have more of an effect on gas transport properties in liquid-like PDMS than in high free volume, glassy PTMSP.

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## Appendix

The activation energy of permeation provides information about the temperature dependence of the permeability coefficient of a penetrant in a polymer. This parameter is usually determined from permeability data measured over a range of temperatures and at a fixed pressure. To assess the contributions of sorption and diffusion to the temperature dependence of the permeation process, the activation energy of permeation is compared with the enthalpy of sorption and the activation energy of diffusion. To make this comparison in a consistent manner, these parameters must also be determined from data at similar conditions of temperature and pressure. Specifically, for the enthalpy of sorption, gas solubility must be determined at the same fixed pressure as the permeability data and over the same range of temperatures. The enthalpy of sorption, thus determined, is an isobaric enthalpy. From eq 14

$$\Delta H_s^P = -R \left( \frac{d \ln S}{d(1/T)} \right)_p \quad (A1)$$

Using eq 4, eq A1 can be rewritten as

$$\Delta H_s^P = -R \left( \frac{d \ln C}{d(1/T)} \right)_p \quad (A2)$$

On the other hand, if the energetics of the sorption process are used to obtain information about differences in interactions of various gas molecules with the polymer matrix, it might be of interest to compare enthalpies of sorption at the same penetrant concentration in the polymer rather than at the same pressure in the gas phase surrounding the polymer. This enthalpy of sorption at constant concentration is the isosteric enthalpy of sorption, which is defined as

$$\Delta H_s^C = -R \left( \frac{d \ln S}{d(1/T)} \right)_C = +R \left( \frac{d \ln p}{d(1/T)} \right)_C \quad (A3)$$

The isobaric and isosteric enthalpies of sorption defined by eqs A2 and A3, respectively, are related, and this relationship is derived below.

To maintain equilibrium between penetrant in a pure gas phase ( $\alpha$ ) and in a sorbed phase ( $\beta$ ) as a result of a change in temperature and/or pressure, the change in penetrant chemical potential,  $\mu$ , in both phases must be equal:<sup>44</sup>

$$d\mu^\alpha = d\mu^\beta \quad (A4)$$

Dividing by the absolute temperature  $T$  and expanding the resulting differentials for  $\mu/T$  in the two phases gives<sup>44</sup>

$$\left( \frac{\partial \mu^\alpha/T}{\partial T} \right)_P dT + \left( \frac{\partial \mu^\alpha/T}{\partial p} \right)_T dp = \left( \frac{\partial \mu^\beta/T}{\partial T} \right)_{P,C} dT + \left( \frac{\partial \mu^\beta/T}{\partial p} \right)_{T,C} dp + \left( \frac{\partial \mu^\beta/T}{\partial C} \right)_{P,T} dC \quad (A5)$$

At constant concentration  $C$ , eq A5 simplifies to

$$\left(\frac{dp}{dT}\right)_C = \frac{\left(\frac{\partial \mu^\alpha/T}{\partial T}\right)_P - \left(\frac{\partial \mu^\beta/T}{\partial T}\right)_{P,C}}{\left(\frac{\partial \mu^\beta/T}{\partial p}\right)_{T,C} - \left(\frac{\partial \mu^\alpha/T}{\partial p}\right)_T} \quad (\text{A6})$$

Substituting for the various partial derivatives<sup>44</sup>

$$\left(\frac{dp}{dT}\right)_C = \frac{-H^\alpha + H^\beta}{T(V^\beta - V^\alpha)} \quad (\text{A7})$$

where  $H^\beta$ ,  $H^\alpha$ ,  $V^\beta$ , and  $V^\alpha$  refer to the partial molar enthalpies and volumes in phases  $\beta$  and  $\alpha$ , respectively. If phase  $\alpha$  is a gas phase and phase  $\beta$  is the sorbed gas in a polymer in equilibrium with phase  $\alpha$ , then typically  $V^\alpha \gg V^\beta$ . Substituting  $(V^\beta - V^\alpha) \approx -zRT/p$  and rearranging<sup>44</sup>

$$\left(\frac{d \ln p}{d(1/T)}\right)_C = \frac{H^\beta - H^\alpha}{zR} \quad (\text{A8})$$

Alternatively, at constant pressure, eq A5 simplifies to

$$\left(\frac{dC}{dT}\right)_P = \frac{H^\beta - H^\alpha}{T\left(\frac{\partial \mu^\beta}{\partial C}\right)_{P,T}} \quad (\text{A9})$$

Now<sup>39</sup>

$$\mu^\beta = \mu_0^\beta + RT \ln a \quad (\text{A10})$$

where  $a$  is the penetrant activity. Combining eqs A9 and A10 and rearranging

$$\left(\frac{d \ln C}{d(1/T)}\right)_P = \frac{H^\beta - H^\alpha}{-R\left(\frac{\partial \ln a}{\partial \ln C}\right)_{P,T}} \quad (\text{A11})$$

From eqs A2, A3, A8, and A11

$$\Delta H_s^C = \frac{1}{z} \left(\frac{\partial \ln a}{\partial \ln C}\right)_{P,T} \Delta H_s^P \quad (\text{A12})$$

Equation A12 provides the relationship between the isobaric and isosteric enthalpies of sorption.

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